Study IR- Raman Spectra properties of Aluminium Phosphide Diamondoids Nanostructures via DFT

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ABSTRACT

Density Functional Theory has been utilized to investigate the electronic and structural characteristics of Aluminium phosphide (AIP). The exchange-correlation potential was calculated using the Generalized Gradient Approximation. The structural, electronic and vibrational features of AIP diamondoids and nanocrystals were investigated using Density Functional Theory at the PBE/6-31(d) level, which included polarization functions. Vibrational modes have been optimized concerning IR intensity, force constants, and lowered masses. In this study there are two components to the vibrational force constant for AIP diamondoids. The first one is distinguished by a reduced mass that is greater than 1 amu and consists primarily of Al-P vibrations that are positioned roughly between 0 and 231 cm⁻¹. The second component has a decreased mass very near to 1 amu and is in the 1228-2400 cm⁻¹ range. It is entirely made up of hydrogen vibrational modes. AIP diamondoids were evaluated with the results of experimental bulk in terms of molecular size-related changes in allocated vibrational frequencies.

KEYWORDS: DFT; AIP; diamondoids nanostructures; IR-Raman spectral.

الخلاصة

استخدمت نظرية دالية الكثافة لدر اسة الخصائص البنيوية والالكترونية لفوسفيد الالمنيوم. وتم استخدام تقريب الميل المعمم للجهد الارتباطي التبادلي. تم حساب الخواص الالكترونية والطيفية للتراكيب البلوية لمويسات فوسفيد الالمنيوم باستخدام نظرية دالية الكثافة صّمن المستوى (BE/6-31(d) المتضمن لدالة الاستقطاب (d). وكذللك تم تعيين الانماط الاهترزازية المتمثلة بثوابت القوة وشدة IR والكتل المختزلة. في هذه الدراسة هناك جز أين لثابت القوة لمويسات فوسفيد الالمنيوم تميز الجزء الاول بقيمة للكتلة المختزلة أكبر من (1amu) والتي تتكون أساساً من اهتزازت AI-P التي يتم وضعها تقريبًا بين (0 cm⁻¹ 0). اما الجزء الثاني له كتلة مُختزلة قريبة جدا من (amu 1) ويقع في مدى (228-2400 cm⁻¹) وهي ضمن انماط اهتزاز الهيدروجين. وتمت مقارنة التغيير المرتبط بالحجم للترددات الاهتزازية المعينة لمويسات فوسفيد الالمنيوم مع القيم التجريبية للأحجام الكبيرة.

INTRODUCTION

At 300 K, AIP can be defined as one of the compound semiconductors with а lattice constant of 0.545 nm. AlP is a semi-conductor with a large indirect band gap. AlP crystallizes in the zinc-blende (zb) structure under normal conditions [1]. It's employed as active material in LEDs and optoelectronic devices, among other things; it's also important in optoelectronic applications in the range of the short wavelength, also for high power, high temperature, and high frequency [2, 3].

Diamondoids are considered nanoscale hydrogen terminated carbon cages [4]. Carbon-based functional nanostructures also referred to as diamondoids, have gotten much attention.

Diamondoids are a class of nanoscale diamond-like carbon nanoclusters that are terminated via hydrogen atoms and are thermodynamically extremely stable [5]. The gap energy, IR spectra, and Raman spectrum of AIP diamondoids have all been investigated by H. Jawad [6]. The research revealed the advantages of using diamantine as an insulating substance to connect drugs to our bodies without interfering with other medications while still delivering medicine to the affected areas. Using the Gauss View and Gaussian 09 software package [7], which refers to an Ab-initio electronic structure program, the optimized structure and vibrational spectra were calculated. The year 09 was used to denote 2009. Many researchers regard Gaussian, a piece of commercial quantum





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chemistry software from Gaussian Incorporated, as the industry standard for molecular modeling and computational chemistry. All of the kev approaches to molecular modeling, including molecular mechanics, ab-initio, semiempirical, and density functional theory (DFT), can be carried out using Gaussian [8]. The ab-initio density functional theory approach, which utilized cluster full geometrical optimization with the complementary Hartree-Fock approach for simulating the optical characteristics and electronic structure of nanocrystals, is a technique of high importance. In addition, this approach is the most resource and time intensive [9].

MATERIALS AND METHODS

Density Functional Theory (DFT) was utilized to perform within Generalized Gradient Approximation (GGA) utilizing Perdew-Burke-Ernzerhof (PBE) exchange-correlation function [10-12]. This study used an even number



of cages with almost equiaxed dimensions, as well as two more small compounds, Al3P3H12 and AlPH6, that are not diamondoids, to demonstrate size effects.

The initial step is to do a geometrical optimization, succeeded by vibrational which is and frequency study. Frequencies are multiplied by a suitable scaling factor (0.986) for assigning the frequencies to the current approach and basis set (PBE/6-31(d)), which must lessen statistical mismatching between the theoretical and the experimental data [13]. This may have compelled us to eliminate a few pf the heavier diamondoids because employing larger basis sets causes computational challenges, which are explained in this paper. the current In work include the geometrically optimized of (Al11P11H28) AlP-tetramantane and AlPoctamantane (Al20P20H42) diamondoids structure molecules as shown in Figure 1a and Figure 1b.



Figure 1. Geometrically optimized a) AlP-tetramantane (Al11P11H28), b) AlP-octamantane (Al20P20H42) diamondoids structure molecules. (by GaussView5.0w).

RESULTS AND DISCUSSION

The use of Raman and IR vibrational spectroscopy for realizing materials is a commonly used in spectroscopy techniques. As a result, it's uncommonly employed to make nanocrystals or determine particle molecular size. The appropriateness of such spectroscopic approaches for creating nanocrystals, AIP diamondoids, and their sizes will be discussed in this paper. Force constants, reduced masses, and vibrations in the Raman spectrum, the reduced mass μ of 2 particles with masses m_b and m_a is calculated as follows [14]:

$$\frac{1}{\mu} = \frac{1}{ma} + \frac{1}{mb} \tag{1}$$

Figures 2 a,b show reduced masses as a function of frequency for both AlP-tetramantane and AlP-octamantane. The vibrations beyond the frequency of 400-438 cm⁻¹ AlP-tetramantane and AlP -octamantane belong to hydrogen-associated vibrations, most of the molecules reduced mass to more than 1 amu. While the reduced molecular masses after 400-438cm⁻¹ are all approximately equal to 1 amu which represents H vibrations.

Figure 3. illustrates the force constants as a function of frequency for AlP-tetramantane and AlP -octamantane. From this figure can be shown

force constants keep on increasing, due to the fact that the vibration frequency is considered to be proportionate to square root related to force constant according to the following equation [15]:

$$\nu = \frac{1}{2\pi} \sqrt{\frac{k}{\mu'}} \tag{2}$$

where k denotes force constant, and $\mu^{\prime} denotes$ the reduced mass.



Figure 2. a) AlP-tetramantane reduced mass as a vibration frequency function and b) AlP-octamantane reduced mass as a vibration frequency function.



Figure 3. a) AlP-tetramantane force constant as a vibration frequency function and b) AlP-octamantane force constant as a vibration frequency function.

Figure 3a and 3b are almost similar in shape and behaviour, these figures might be divided to 2 parts. The first one starts from about 0cm⁻¹ and ends at not more than 231-390cm⁻¹, while the second one begins at 231 -390cm⁻¹ and ends at about 2400cm⁻¹. The second part in Figure 3 can easily explains the parabola shape, because reduced mass is fairly constant and equals 1. It can be seen in the second part of Figure 2 and Figure 3 have similar shapes and values compared to the first part for the same figures.

High reduced mass mode (HRMM) and high force constant mode (HFCM) with different AlP diamondoids are shown in Table 1. Tetramantane diamondoids' (HRMM) and (HFCM) structures are smaller than those of octamantane.

Table 1. High Reduced Mass Mode (HRMM) and High Force Constant Mode (HFCM).

AlP diamondoids	High Reduced Mass Mode (HRMM) cm ⁻¹	High Force Constant Mode (HFCM) cm ⁻¹
Tetramantane (Al11P11H28)	400.07	231.81
Octamantane (Al20P20H42)	438.59	390.11

Figure 4 shows the AIP diamondoid size variations of the 4 different sizes. Longitudinal Optical-Highest reduced Mass Mode (LO-HMM), Radial Breathing Mode (RBM), and Al-2H and P-2H asymmetrical vibrations are among such vibrations. Through the observation of the atomic displacement regarding such vibration, RBM can be achieved. In this mode, the majority of atoms move in phase along the molecule's radius When reaching bulk sizes, the [16]. vibration frequency must converge to 0 cm^{-1} [17]. The LO-HRMM, on the other hand, converges to 441cm⁻¹[18]. In the AlPH6 molecule, the 2 modes meet at 240 cm⁻¹. Therefore, RBM is blue shifted as it moves from bulk to nanoscale and molecular limits, whereas LO-HRMM is red shifted. HRMM and RBM, on the other hand, vary in size and might be utilized as size directors for those diamondoids.

Figures 5a and 5b illustrate the IR spectra as a function of frequency for AlP-tetramantane and AlP-octamantane diamondoids structure molecules respectively.



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AlP-octamantane



structure



Figure 4. Molecular size effects on various AlP diamondiod molecular vibration values, HRMM, RBM, and asymmetric hydrogen surface vibration modes are seen. The dashed line denotes experimental bulk AlP LO vibration frequency [18].



Figure 5. IR spectra as a function of frequency for a) AlP-tetramantane b) AlP-octamantane diamondoids structure molecules.

There are 2 regions based on the characteristics of vibration or gap separation. In the first region, the results indicate good matching to experimentally acquired lines. The (0-570) cm⁻¹ region is characterized by pure AIP. Notice a shift in the intensity maxima toward the right side of the infrared vibrational frequencies which is indicative of the presence H surface atoms. The second region includes the 1066–1870 cm⁻¹ modes in the Al-H and P-H vibrational regions.

Figures 6a and 6b show the Raman spectra as a function of frequency for AlP-tetramantane and



diamondoids

Figure 6. Raman spectra as a function of frequency for (a) AlP-tetramantane (b) AlP-octamantane diamondoids structure molecules.

Raman spectra for AlP diamondoid structure are between the 1880-2385 cm⁻¹ modes in the Al-H and P-H vibrational region including stretching (symmetrical and asymmetrical) and bending vibration mode (wagging, rocking, scissoring, and twisting). The corresponding Raman spectra share almost identical characteristics as the IR spectra, however, with different selection rules. The chance of absorbing a specific wave length photon determines the peak intensities in the Raman spectrum. These results have been discovered to be in consistency with other theoretical values in reference [19].

CONCLUSIONS

The AIP vibration and characteristics estimated with the use of the DFT approach are found to be in good agreement with experimental bulk limits in the current study. Various spectroscopic and structural features were investigated using molecular calculations. The theoretically sensitive size variations of AIP diamondoids have been demonstrated using IR Raman spectroscopic techniques. There are two parts to the vibrational force constant for AIP diamondoids. The first consists primarily of AlP vibrations with a decreased mass of more than 1 amu and a range of around 0-231 cm⁻¹. The second part, which is in the range 1228-2400cm⁻¹ and has a decreased mass of very near to 1 amu, is made up entirely of hydrogen vibrational modes. When moving from the bulk to nanoscale or molecular, RBM experiences a blue shift; while when moving from the bulk to molecular or nanoscale, LO-HRMM experiences red shift.

Disclosure and conflict of interest: The authors declare that they have no conflicts of interest.

REFERENCES

- Mujica A, Rodríguez-Hernández P, Radescu S, Needs RJ, Muñoz A. AlX (X = As, P, Sb) Compounds under Pressure. Phys Stat Solidi (b); 211, (1999) pp 39-43. <u>https://doi.org/10.1002/(SICI)1521-</u> 3951(199901)211:1<39::AID-PSSB39>3.0.CO;2-N
- I. Vurgaftmana & J. R. Meyer, "Band Parameters for III-V Compound Semiconductors and their Alloys", Journal of Applied Physics, 89; (2001) 5815. <u>https://doi.org/10.1063/1.1368156</u>
- [3] V. B. Sheeba & C. N. Louis, "Band Structure, Metallization and Superconducting Transition of Group III-V Semiconductors AlP and AlAs under High Pressure", Chemical and Materials Engineering 1(4): (2013) 132-140. https://doi.org/10.13189/cme.2013.010405
- [4] Dahl J E, Liu S G and Carlson," Isolation and Structure of Higher Diamondoids, Nanometer-Sized Diamond Molecules ".R M K Science, 299, (2003) 96-99. <u>https://doi.org/10.1126/science.1078239</u>
- [5] Lifshitz Y, Köhler T, Frauenheim T, Guzmann I, Hoffman A, Zhang R Q, Zhou X T and Lee S T," The Mechanism of Diamond Nucleation from Energetic Species ",.Science 297, (2002) pp1531- 1533. <u>https://doi.org/10.1126/science.1074551</u>
- [6] Huda M. Jawad, "Building nanoparticles using diamondoids structures as a carrier of the drug: Density functional theory study", Journal of Garmian University, https://doi.org/10.24271/garmian.136, (2017).

https://doi.org/10.24271/garmian.136

[7] Frisch, M.J., G.W. Trucks and H.B. Schlegel, Gaussian 09, Revision A.02, Gaussian, Inc., PA, Wallingford CT., (2009).

- [8] Yousif Abid Al Shaabani, Bahjat B. Kadhimb, Fikrat H. Jasim, Structural and Electronic Properties of Theophylline- InPDiamantane Drug Carrier, American Scientific Research Journal for Engineering, Technology, and Sciences, Vol. 19, No 1, (2016) pp 130-141.
- [9] M.T. HusseinÅ, M. A. Abdulsattar and H.A. Hameed, "Electronic and vibrational spectroscopic properties of GaAs diamondoids using Density Functional Theory", International Journal of Advance Industrial Engineering, .2, No.4 (2014). <u>https://doi.org/10.1063/1.4904031</u>
- [10] Hohenberg P and Kohn W ," Inhomogeneous Electron Gas" ,Phys. Rev. 136 B864 (1964). https://doi.org/10.1103/PhysRev.136.B864
- [11] Kohn W and Sham , L J Phys. Rev. 140 A1133 (1965). <u>https://doi.org/10.1103/PhysRev.140.A1133</u>
- Perdew J P, Burke K and Ernzerhof M," Generalized Gradient Approximation Made Simple", Phys., Rev. Lett. 77 No.18, (1996). https://doi.org/10.1103/PhysRevLett.77.3865
- [13] Russell D. Johnson III, "NIST Computational Chemistry Comparison and Benchmark Database ",Reference Database Number 101 Release, (1999).
- [14] Mudar A. Abdulsattar, Size Dependence of Si Nanocrystals Infrared Spectra: ADensity Functional Theory Study, Silicon, 5: (2013) pp 229-237. <u>https://doi.org/10.1007/s12633-013-9152-4</u>
- [15] Fatima A. Hasan, Mohammed T. Hussein," Study of some electronic and spectroscopic properties of ZnO nanostructers by density functional theory", Materials Today: Proceedings 42 (2021) 2638-2644. <u>https://doi.org/10.1016/j.matpr.2020.12.593</u>
- [16] Abdulsattar M A," Size variation of infrared vibrational spectra from molecules to hydrogenated diamond nanocrystals: a density functional theory study". Beilstein J. Nanotechnol. 4: No.262. (2013) 262-268 <u>https://doi.org/10.3762/bjnano.4.28</u>
- [17] Kittel C , "Introduction to Solid State Physics", eighth ed. New York: Wiley, (2005).
- [18] Otfried Madelung, "Semiconductors: Data Handbook", Published by Springer Berlin Heidelberg, ISBN: 978-3-642-62332-5, (2004).
- [19] Hayder M., Mudar A. Adudlsattar, "Strutural and spectroscopic properties of AlP Diamondoids: A DFT study", J. Eng. Applied Sci., 13(12), (2018) pp 4381-4386.

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